

## Energy harvesting of a pendulum vibration absorber

*Abstract.* This paper presents the numerical results of a pendulum vibration absorber. The system composed of the pendulum attached to the Duffing's oscillator. The harvester is mounted in the pendulum pivot structure and consists of a cylindrical permanent magnet mounted on rotor and of four windings fixed to the housing as a stator. The obtained results show the influence of system and electrical parameters on value of induced voltage for vibration absorption conditions.

**Streszczenie.** W pracy przedstawiono badania numeryczne wahadłowego tłumika drgań. Układ składa się z wahadła zamocowanego na oscylatorze Duffinga. Urządzenie indukujące prąd zamocowane jest w zawieszeniu wahadła i składa się z cylindrycznego magnesu zamontowanego na rotorze i czterech uzwojeń zamontowanych na stojanie. Otrzymane wyniki pokazują indukowane napięcie w warunkach dynamicznej eliminacji drgań. (**Odzyskiwanie energii w wahadłowym tłumiku drgań**).

**Keywords:** wahadło, odzyskiwanie energii, rezonans, indukcja, voltage

**Słowa kluczowe:** pendulum, energy harvester, resonance, induction, napięcie.

### Introduction

Reduction of the vibration is a serious concern for many practical applications. For example, tall buildings and long-span bridges are subjected to huge dynamic loadings from the winds, earthquakes, water waves, traffics and human motions. The large vibration amplitude can damage the structures or the secondary components, or cause discomfort to its human occupants [1]. Among of many methods to reduction of vibration, the pendulum mass damper (PMD) has been proved to be a very simple and effective vibration suppression device, with many practical implementations on tall buildings, such as Taipei 101 in Taipei, Citi Group in New York, and many others [2].

Among the nonlinear systems, we can mention a special class of models which consists of at least two subsystems. If the subsystems are coupled by inertial terms, then periodic vibrations generated by one subsystem become an excitation source for the other. This phenomenon performs the so called autoparametric vibrations. Recent developments of absorption devices include the autoparametric vibration absorbers. The idea lies in attaching the absorber to the primary system in such a manner that it experiences a parametric base excitation, and therefore, the absorber frequency is tuned around one-half of the troublesome frequency value [3]. The governing system of equations of motion has quadratic nonlinearities, whose influence is activated by this tuning, enabling a parametric resonance to be excited and the absorber to respond in the desired frequency range. This device provides vibration suppression of the primary system response.

We propose to apply harvesting energy device in the autoparametric pendulum vibration absorber (APVA), this solution allows simultaneous generate electrical energy and maintain the vibration absorption effect. Autoparametric vibration systems have an interesting dynamic that results from at least two non-linear subsystems coupled to interact in a way where one of them transfers the exogenous perturbation energy to the other (vibration absorption effect). Thus, the primary system can be externally excited by some harmonic force and, when it is connected to the secondary system (absorber), it can be verified the so called parametric excitation, that is, a mechanism that transfers the exogenous energy to its absorber. The autoparametric interaction has been analyzed in the literature: (Haxton and Barr [4], Cartmell and Lawson [5], Korenev and Reznikov [6], Nayfeh and Mook [7], Warminski

and Kęcik [8]). There are many vibrating structures and machinery, where the pendulum type absorbers are used. Energy harvesting has been an active research area and systems refer to devices that capture and transform energy into electricity. Usually, the energy could be the kinetic energy of moving or vibrating structures. Vibration-based energy harvesters can be categorized into three types, namely electromagnetic, piezoelectric, and electrostatic, depending on the medium of the transducer. The transducers require input vibration with high frequency to generate electricity. Civil structures exhibit usually low frequency structural vibrations, which make it difficult for an energy harvester to extract energy. Therefore, tuned in such systems is very important. Vibrational energy harvesters achieve their highest output power when they are working nearly resonance regions (because of the highest amplitude of vibration).

Research in the area of piezoelectric energy harvesting involves understanding the mechanics of vibrating structures, the fundamental electrical circuit theory and the constitutive behavior of piezoelectric materials. The modeling problem of vibration energy harvesting using piezoelectric transduction is to estimate the voltage output across the resistive load in terms of the base motion input. The voltage output can then be used to calculate the power delivered to the given electrical load.

This paper shows the numerical study of an autoparametric pendulum absorber with mounted harvesting system. The study is done nearly region, where the dynamic vibration absorption effect is satisfied. Moreover, the effective electricity regarding different system and electrical parameter is discussed.

### Model of an autoparametric pendulum harvester

The main idea is to generate the electricity by convert the kinetic energy (the rotation or swings of the pendulum) into the voltage. The model of an autoparametric pendulum harvester in Fig. 1 is shown. This model is composed of the harmonic excited oscillator (the mass  $m_1$ ) and attached pendulum (the masses  $m_2$  and  $m_3$ ). The stiffness of springs are denote:  $k$  and  $k_2$  - linear characteristics,  $k_1$  - nonlinear characteristic. In the pendulum pivot, harvesting device is mounted. This device consist of a rotor a cylindrical permanent magnet and a stator. The device is connected directly to the electrical load,  $L_d$ . The scheme of harvesting device mounted in the pendulum pivot in Fig. 2 is shown.

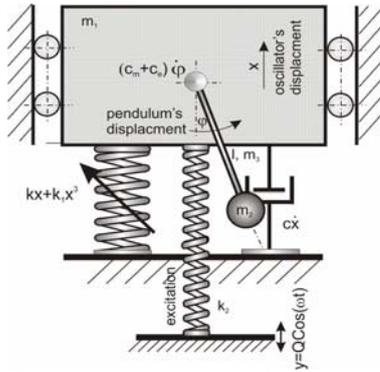


Fig. 1. Model of the autoparametric harvester.

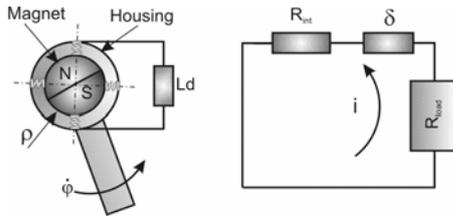


Fig. 2. Model of the energy converter and equivalent circuit.

The equations of motion are derived by Lagrange's equations of the second kind. The detailed derivation of dimensional equations can be found in [9]. They are written as:

$$(1) \quad (m_1 + m_2 + m_3) \ddot{x} + c\dot{x} + (k + k_2)x + k_1x^3 + \left(m_2 + \frac{m_3}{2}\right) l (\ddot{\phi} \sin \phi + \dot{\phi}^2 \cos \phi) = k_2 Q \cos \omega t,$$

$$(2) \quad \left(m_2 + \frac{m_3}{3}\right) l^2 \ddot{\phi} + (c_e + c_m) \dot{\phi} + \left(m_2 + \frac{m_3}{2}\right) l (\ddot{x} + g) \sin \phi = 0.$$

The first, Eq. (1) represents the oscillator's motion while the second Eq. (2) denotes motion of the pendulum. If the pendulum does not move, it becomes a part of the oscillator mass, only. The  $c_m$  is the damping coefficient used to describe the mechanical damping,  $c_e$  denotes electrical damping in the electrical devices mounted in the pendulum pivot and  $c$  is a viscous damping coefficient in damper.

To calculate the electrical damping coefficient, the equation for the electrical circuit is used by applying Kirchhoff's law to the electrical circuit

$$(3) \quad i(R_L + R_I) - \delta \dot{\phi}(\tau) \rho = 0,$$

where:  $\delta$  is the transduction factor,  $\rho$  is the radius of housing

on which the coils are mounted,  $R_L$  is the resistance of the external load,  $R_I$  is the internal resistance. The  $\delta$  equals the magnetic flux gradient. This term is commonly depends on: number of coil turns, average magnetic field strength, and the coil length. The voltage can be estimated based on Ohm's law

$$(4) \quad U = iR_L = \delta \frac{R_L \dot{\phi}(t) \rho}{R_L + R_I}.$$

Then the electrical damping coefficient defined as

$$(5) \quad c_e = \frac{\delta^2 \rho}{R_L + R_I}.$$

The value of this coefficient depends on electrical properties and construction of the harvester device.

Introducing dimensionless time  $\tau = \omega_0 t$  where  $\omega_0$  is the natural frequency of the linear oscillator (the nonlinear part comes from the nonlinear spring is neglected), and then the

dimensionless coordinates  $X = x/x_{st}$  and  $\varphi = \phi$  where  $x_{st}$  is the static displacement of the linear oscillator we obtain the equations (1) and (2) in dimensionless form:

$$(6) \quad \ddot{X} + \alpha_1 \dot{X} + X + \gamma X^3 + \mu \lambda (\dot{\varphi} \sin \varphi + \dot{\varphi}^2 \cos \varphi) = q \cos \vartheta \tau, \\ \ddot{\varphi} + \alpha_2 \dot{\varphi} + \beta \varphi + \lambda (\ddot{X} + 1) \sin \varphi = 0.$$

The dimensionless parameters take the following definitions:

$$\alpha_1 = \frac{c}{(m_1 + m_2 + m_3) \omega_0}, \alpha_2 = \frac{c_m}{\left(m_2 + \frac{m_3}{3}\right) l \omega_0}, \\ \beta = \frac{\delta^2 \rho}{(R_L + R_I) \left(m_2 + \frac{m_3}{3}\right) l \omega_0}, q = \frac{k_2 Q}{(k + k_2) x_{st}}, \\ (7) \quad \omega_0 = \sqrt{\frac{k + k_2}{(m_1 + m_2 + m_3)}}, x_{st} = \frac{(m_1 + m_2 + m_3) g}{(k + k_2)}, \gamma = \frac{k_1}{k + k_2} x_{st}^2, \\ \mu = \frac{\left(m_2 + \frac{1}{2} m_p\right) l^2}{(m_1 + m_2 + m_3) x_{st}^2}, \lambda = \frac{\left(m_2 + \frac{1}{2} m_3\right) x_{st}}{\left(m_2 + \frac{1}{3} m_3\right) l}, \vartheta = \frac{\omega}{\omega_0}.$$

The motion of the dynamic pendulum absorber is described by two generalized coordinates, namely the displacement of the oscillator in the vertical direction  $X$ , and the angle of the pendulum rotation  $\varphi$ .

### Harvesting energy and absorption effect

The autoparametric vibration pendulum absorber (AVPA) is designed to absorb energy from the primary system (main mass). This absorption effect is efficient only in the limited band of vibration frequencies of the main system. Therefore we analyze dynamics both systems i.e. the pendulum and the oscillator motion. The system parameters are:  $\alpha_1=0.1$ ,  $\alpha_2=0.002$ ,  $\mu=6$ ,  $\lambda=0.3$ ,  $q=0.2$ ,  $\gamma=0$ . The electrical parameters are:  $R_I=188\Omega$ ,  $R_L=10^6\Omega$ ,  $\delta=7.752\text{Vs/m}$ ,  $\rho=0.005\text{m}$ . The simulation time was taken for two thousand periods where first thousands of them were cut as a transient time. The calculations have been done in Matlab environment by means of Runge–Kutta method using solver ode45. The results of this work based of the voltage root mean square (quadratic mean) which derivatives  $RMS_V = \langle V \rangle / \sqrt{2}$ . This is a statistical measure of the magnitude of a varying quantity. It is especially useful when variates are positive and negative.

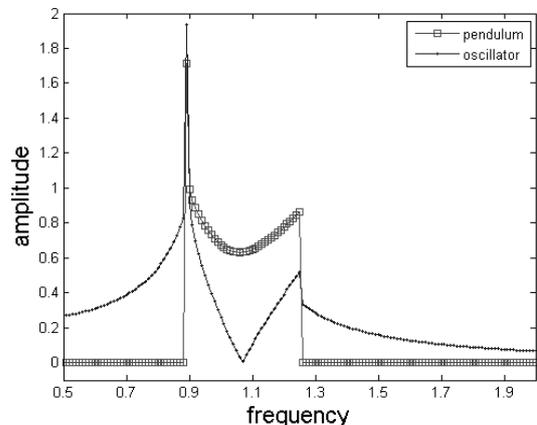


Fig. 3. The resonance curves for the pendulum (blue) and the oscillator (red).

The resonance curves in Fig. 3 are presented. The line (·) and (◦) denote the pendulum and the oscillator resonances, respectively. Near the frequency  $\vartheta=1.1$ , the dynamical vibration absorption effect is clearly visible, the amplitude of oscillator is reduced to zero, the full energy takes the pendulum.

The amplitude of excitation ( $q$ ) is a crucial parameter which influence on the dynamics of an autoparametric system. The generating electricity depends on this parameter, also. This situation in Fig. 4 is presented. Increase amplitude of excitation causes, that resonance region is expanded and induced electricity growing. Unfortunately, this parameter may cause reduction in efficiency of absorption effect (the chaotic motions or the rotations of the pendulum is possible [9]). The maximal value of root mean square of voltage equals about 1.5.

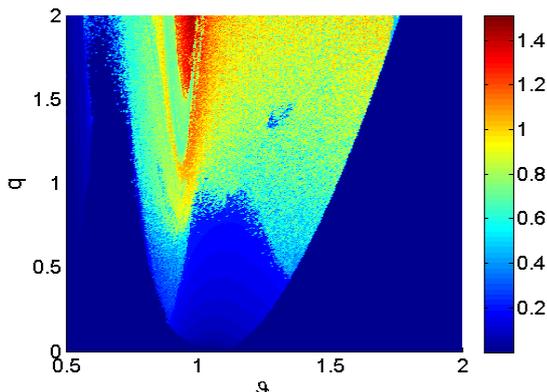


Fig.4. Influence of amplitude excitation ( $q$ ) versus frequency of excitation ( $\vartheta$ ) on the RMS<sub>v</sub> voltage near the resonance region.

The nonlinearity of a supporting spring ( $\gamma$ ) slightly influences location of the resonance and absorption regions. The positive characteristic of spring shifted the resonance region in right, the negative characteristic in left side (Fig. 5). From, the practical point of view, very important is that nonlinearity of the spring practically doesn't decrease the electricity generated, but the dangerous region can be move away.

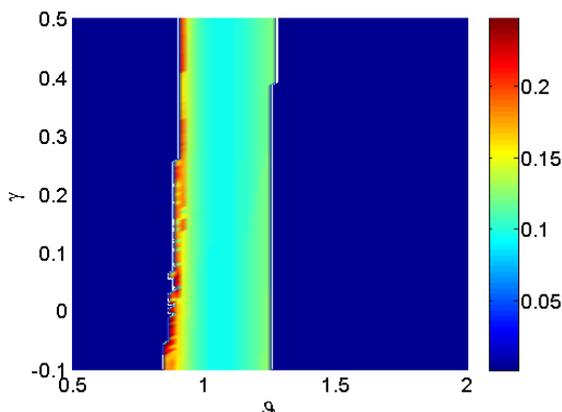


Fig.5. Influence of spring's nonlinearity ( $\gamma$ ) versus frequency of excitation ( $\vartheta$ ) on the RMS<sub>v</sub> voltage near the resonance region.

The dynamics of the system (and of course induced electricity) strongly depend on the values of the nonlinear terms which couple the main structure (the oscillator) and absorber (the pendulum). The response of the system is very sensitive to changes to the parameters  $\lambda$  and  $\mu$ . Parameter  $\lambda$  and  $\mu$  couples the oscillator and pendulum motion. Parameter  $\mu$  only appears in the oscillator equation of motion and plays the role of a gain in the product  $\mu\lambda$  in

the oscillator equation (first in Eq. 6). In a real system, the length of the pendulum and its mass moment of inertia are responsible for these parameters values (see eq. (7)). The influence of both parameters  $\lambda$  and  $\mu$  on the value of RMS<sub>v</sub> voltage in Fig. 6 is presented. By combination of these parameters, for the analyzed range, we can increase several times in the RMS<sub>v</sub> voltage. Additionally, the region where the electricity is induced is clearly visible (the pendulum executes motion). For,  $\lambda < 0.15$ , the pendulum stopped and the electricity is not generated.

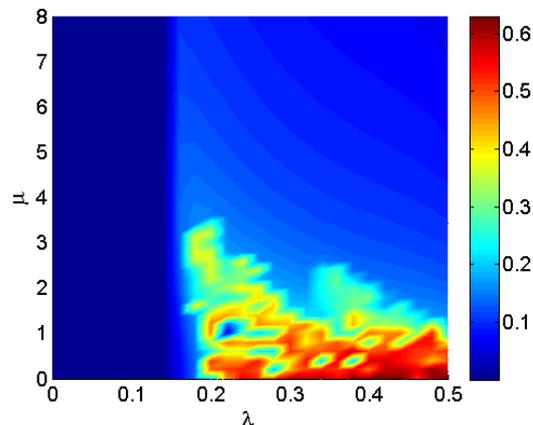


Fig.6. Influence of pendulum's parameters ( $\mu$  versus  $\lambda$ ) on the RMS<sub>v</sub> voltage, for  $\vartheta=1.05$ .

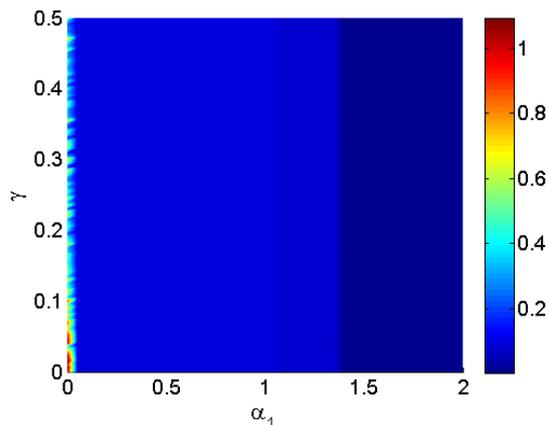


Fig.7. Influence of suspension's parameters ( $\gamma$  versus  $\alpha_1$ ) on the RMS<sub>v</sub> voltage, for  $\vartheta=1.05$ .

The suspension of an autoparametric vibration absorber can be used for a control of dynamics of absorber in difficult situations. Interestingly, that induced electricity practically is independent on the spring stiffness ( $\gamma$ ), especially for larger oscillator damping. This results is similar with Fig 5, where nonlinearity of spring slightly shifted the resonance frequency. Using damping in the suspension of an autoparametric vibration system we can adjust the value of electricity induced. In, Fig. 7, we can see the influence of oscillator damping versus nonlinearity of spring. Analyzing, the results, we can conclude, that control of electricity is easy by coefficient  $\alpha_1$  (for example: changing damping in viscous damper by hydraulic valve or current in magnetorheological damper). These region is a wide, therefore easily adjust the system parameters to the required voltage.

Of course, the parameters of harvester device system is also important. The dimensionless electrical damping ( $\beta$ ) depends the radius of housing ( $\rho$ ), magnetic flux gradient ( $\delta$ ) and the resistance of the external and internal load ( $R_L$  and  $R_i$ ). The influence of the electrical damping ( $\beta$ ) causes

reduction both to the right and the left hand sides of the resonance region, which essentially decreases the effectiveness of the pendulum as a dynamical damper (Fig. 8). Interestingly, that maximal value of generate electricity exists for small region, only. Similar effect appears for mechanical damping of the pendulum ( $\alpha_2$ ), because the electrical damping is a part of total damping of the pendulum ( $\alpha_2 + \beta$ ). The total damping is an important parameter, for dynamic absorption of vibrations. A slightly difference in coefficients ( $\alpha_2$  or  $\beta$ ) influences effectiveness absorption effect. Increase damping in the pendulum causes decrease of efficiency of absorption effect [7]. For the damping in the harvester device about  $\beta \approx 0.5$ , the pendulum stopped and electricity is not induced. For our parameters the electrical damping coefficient is too small of the order  $3 \cdot 10^{-6}$ .

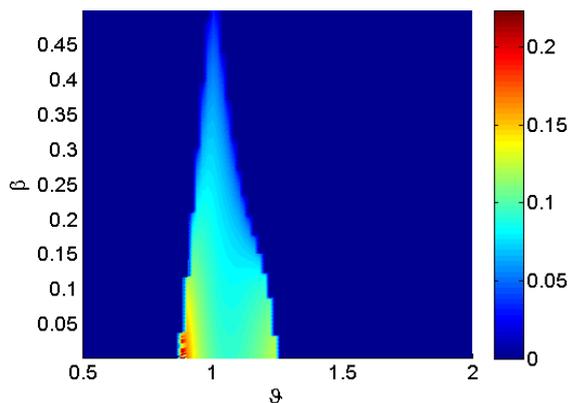


Fig.8. Influence of electrical damping coefficient ( $\beta$ ) on the RMS<sub>v</sub> voltage near the resonance region.

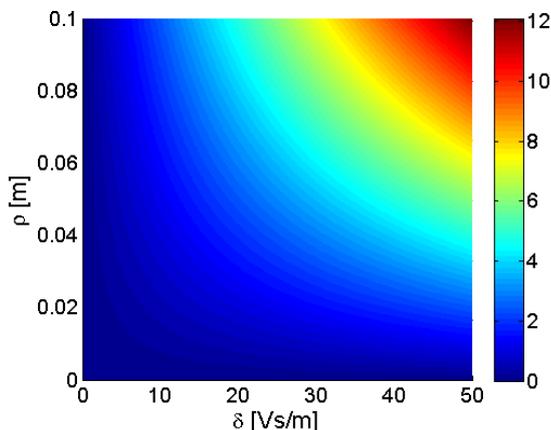


Fig.9. Influence of harvester device parameters ( $\rho$  versus  $\delta$ ) on the RMS<sub>v</sub> voltage for the absorption effect,  $q=1.05$ .

Radius of housing versus magnetic flux gradient, obtained for linear spring  $\gamma=0$  is presented in Fig. 9. These parameters significantly influences electricity. For, example, for  $\rho=0.06\text{m}$ , the voltage root mean square can achieve about 7V (for  $\rho=50\text{Vs/m}$ ).

It is worth noting that most of the numerical results presented in this work (without installed harvesting device) were verified analytically and experimentally on the special laboratory system [8,10].

## Conclusions and final remarks

The vibration absorption systems are device dedicated to reduce dangerous vibration. This work proposes the application of energy harvesting device, mounted in the pivot of the pendulum vibration absorber. Because of the small electrical damping (see eq. 7) the effectiveness of absorption effect is practically unchanged. The excessive increase in damping of the pendulum causes decaying a big reduction of the resonant curve (Fig. 8), an unwanted effect from the vibration dynamical absorption point of view.

To control of dynamic and voltage induced can be realized by proper choice of the pendulum parameters ( $\mu$  and  $\lambda$ ), but these parameters cannot be modified online, as the pendulum operates. Therefore, another solution is proposed, based on control the damping and the stiffness of a suspension of an autoparametric vibration pendulum absorber. The amplitude of excitation ( $q$ ) influences on absorption effect, therefore in this type of absorbers required to apply the control with feedback.

In future, the harvesting device will mounted in the laboratory rig, and these results will be experimentally verified.

## Acknowledgments

The work is financed by Grant no. 0234/IP2/2011/71 from the Polish Ministry of Science and Higher Education in years 2012-2014.

## REFERENCES

- [1] Kareem A., Kijewski T., Tamura Y., Mitigation of Motions of Tall Buildings with Specific Examples of Recent Applications, *Wind and Structures*, 2(1999), 3, 201-251.
- [2] Bungale S.T. *Reinforced Concrete Design of Tall Buildings*. CRC Press, 2010.
- [3] Regis V. Tuning Methodology of Nonlinear Vibration Absorbers Coupled to Nonlinear Mechanical Systems. PhD Thesis, 2010, University of Liege.
- [4] Haxton R.S., Barr D. S., The autoparametric vibration absorber," *ASME Journal of Engineering for Industry*, 94(1972), 119-125.
- [5] Cartmell M.P., Lawson J.W., Performance enhancement of an autoparametric vibration absorber by means of computer control, *Journal of Sound and Vibration*, 177(1994), 173-195.
- [6] Korenev B.G., Reznikov, L.M. *Dynamic Vibration Absorbers: Theory and Technical Applications*, Chichester: John Wiley & Sons, 1993, UK.
- [7] Nayfeh A.H., Mook D.T., *Nonlinear Oscillations*. New-York: Wiley-Interscience, 1995.
- [8] Warminski J., Kecik K. Autoparametric vibrations of a nonlinear system with a pendulum and magnetorheological damping. *Nonlinear Dynamics Phenomena in Mechanics*. Eds. J. Warminski, S. Lenci, M. P. Cartmell, G. Rega and M. Wiercigroch, Springer, 2012.
- [9] Warminski J., Kecik K. Instabilities in the main parametric resonance area of mechanical system with a pendulum. *Journal of Sound Vibration*, 332 (2009), 612-628.
- [10] Kęcik K. Zastosowanie tłumika magnetoreologicznego do sterowania drganiami w układzie mechanicznym z wahadłem. *Przegląd Elektrotechniczny (Electrical Review)*, 2(2012), 223-226.

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